

Strongly-interacting phases in heavy-ion collisions and neutron stars

Ayon Mukherjee

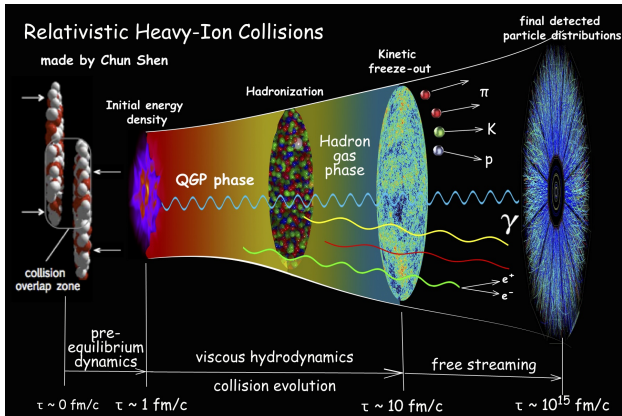
with

M. Csanád, V. Dexheimer, T. Galatyuk, A. Jakovác, S. Lökös, R. Rapp, S. Schramm,
F. Seck, J. Steinheimer, H. Stöcker, J. Stroth, S. Tripathy, M. Wiest

&

data from the STAR Collaboration

April 7th, 2023



[Chun Shen, The Ohio State University]

- HIC's to probe QCD phase structure & CEP
- Pheno. model \rightarrow EoS \rightarrow Hydro. sim. of eqibm. stage \rightarrow phase structure
- Exp.-data analysis \rightarrow femtoscopy \rightarrow source-geometry \rightarrow phase structure

PART I: PHENOMENOLOGY

Ideal, relativistic hydrodynamics

- Macroscopic description of ideal fluid requires conserved quantities
- Ideal fluid: a continuous system of infinitesimal volume elements, each of which are assumed to be very close to thermodynamic equilibrium
- Conservation laws: $\nabla_{\mu} T^{\mu\nu}_{(0)} = 0$, $\partial_{\mu} N^{\mu}_{(0)} = 0$
- Fields: ε, P, n and u^{μ} correspond to 6 degrees-of-freedom
- Equations of motion:

$$\begin{aligned} D\varepsilon + (\varepsilon + P)\theta_{\mu} u^{\mu} &= 0 \\ (\varepsilon + P)Du^{\alpha} + c_s^2\theta^{\alpha}\varepsilon &= 0 \\ Dn + n\partial_{\mu} u^{\mu} &= 0 \\ c_s^2(\varepsilon) - \frac{\partial P(\varepsilon)}{\partial \varepsilon} &= 0 \end{aligned}$$

- Equation-of-State: $P \equiv P(n, \varepsilon)$ from thermodynamic model based on microscopic theory of strong interactions

The $Q\chi P$ model

- Description:

- Flavour SU(3) extension of non-linear representation of σ - ω model
- Grand-canonical, thermodynamic model
- Effective mass of baryons:

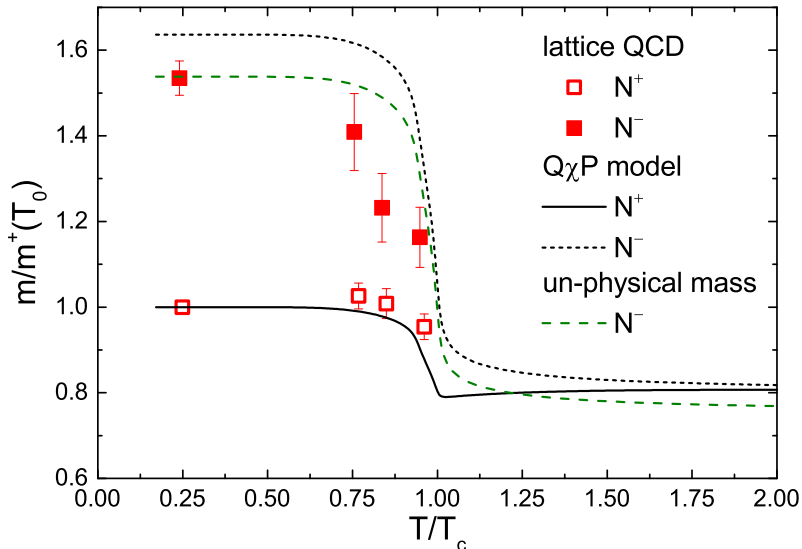
$$m_{i\pm}^* = \sqrt{\left[(g_{\sigma i}^{(1)}\sigma + g_{\zeta i}^{(1)}\zeta)^2 + (m_0 + n_s m_s)^2 \right]} \pm g_{\sigma i}^{(2)}\sigma \pm g_{\zeta i}^{(2)}\zeta$$

- Order-parameter for chiral transitions: $\sigma = \langle \bar{\psi}\psi \rangle$
- Order-parameter for deconfinement: Polyakov loop, ϕ
- Hadrons removed, post deconfinement, with excluded-volumes

- Objectives:

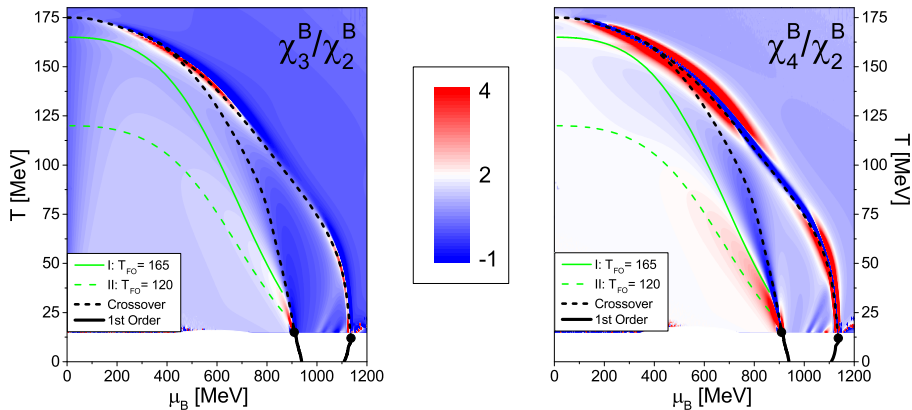
- Qualitative agreement with lattice-QCD predictions
- χ_B^n near both phase transitions & critical end-points – LG & FOPT
- Phase-diagram at $|\mu_S| \geq 0$; with $\mu_B \geq 0$ & $\mu_I = 0$
- Properties of ground-state nuclear-matter & neutron-star-matter
- Quantitative agreement with UrQMD-simulation outcomes
- Hydrodynamic simulations of HIC's with model EoS's
- Quantitative agreement with coarse-grain-transport results & HADES data

Lattice-QCD comparison



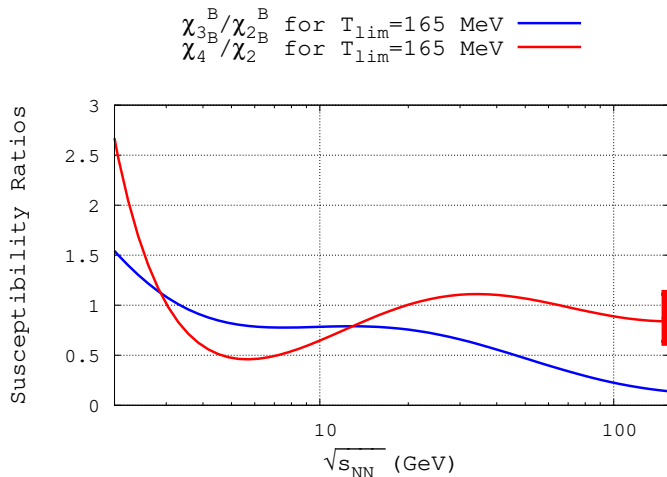
- Mass-degeneracy of parity-doublets at $T > T_{CEP}$ consistent with IQCD

Phase diagrams



- Cumulant-ratios increase near CEP's & near crossover-merger at low μ_B
- Enhancement more pronounce for χ_4^B/χ_2^B

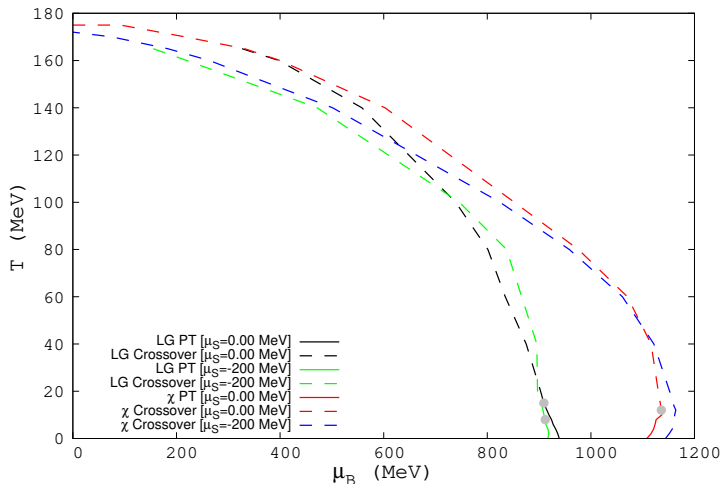
Susceptibilities



[A. Mukherjee, J. Steinheimer & S. Schramm; Phys. Rev. C 96 (2017) 2, 025205]

- Cumulant-ratios deviate considerably from unity

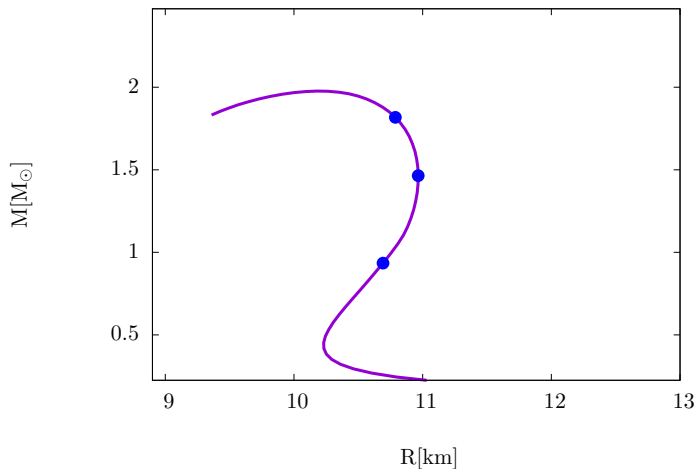
Strangeness



[A. Mukherjee, A. Bhattacharyya & S. Schramm; Phys. Lett. B 797 (2019) 134899]

- Chiral-FOPT disappears for $|\mu_S| \geq 175$ MeV due to hyperon-domination

Astrophysical benchmarks



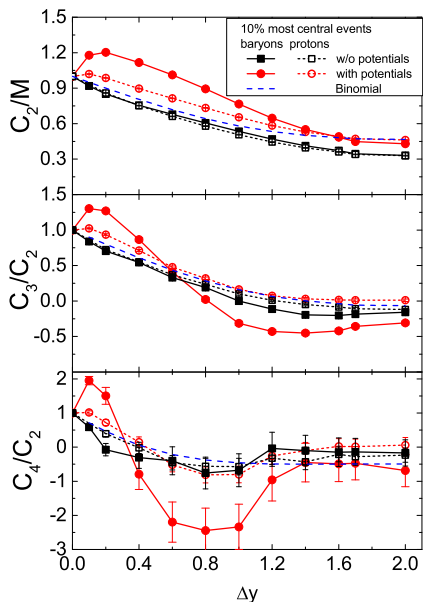
[A. Mukherjee, J. Steinheimer, S. Schramm & V. Dexheimer; *Astron. Astrophys.* 608 (2017) A110]

- Model-EoS with TOV to generate M - R diagram for neutron stars
- Maximum mass & radius in agreement with observations

'Numbers' speak louder than words!

- Ground-state nuclear-matter $\kappa = 267.12$ MeV
- Saturation density (ρ_0) = 0.142 fm^{-3}
- Binding energy (E/A) or, energy-density per baryon (ε/ρ_B) = -16 MeV
- Symmetry energy: $S = \frac{1}{8} \left[\frac{d^2(\varepsilon/\rho_B)}{d(l_3/B)^2} \right]_{\rho_B=\rho_0} = 30.02$ MeV
- Slope parameter: $L = 3\rho_0 \left[\frac{dS}{d\rho_B} \right]_{\rho_B=\rho_0} = 56.86$ MeV
- Maximum star mass: $M_{\text{max}} = 1.98 M_{\odot}$
- Maximum star radius: $R_{\text{max}} = 10.25$ km
- Canonical star mass: $M_c = 1.4 M_{\odot}$
- Canonical star radius: $R_c = 11.10$ km

UrQMD with nuclear potentials



- Enhanced cumulant-ratios with nuclear potentials, for $\Delta y < 0.3$
- For $\Delta y > 0.3$, all C_n suppressed, due to baryon-number conservation
- Enhancement smaller for net- N_p than net- N_B , due to random exchange of isospin with neutrons & pions
- Cascade mode agrees with simple binomial distr. for net- N_B

[J. Steinheimer et al; Phys. Lett. B 785 (2018) 40–45]

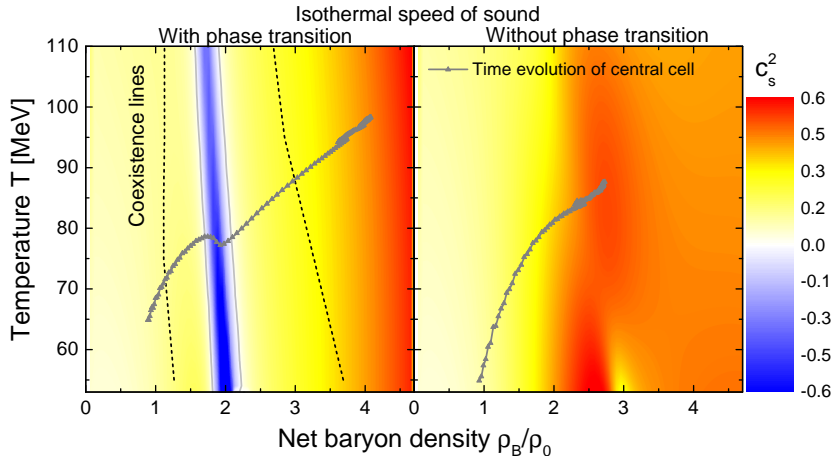
Dileptons

- Dileptons: effective probes for the evolution of the fireball; on account of electro-weak interactions being unlikely at strong-interaction timescales
- Invariant-mass spectrum of dileptons obtained from emissivity

$$\epsilon = -\frac{\alpha_{EM}^2}{\pi^3} \frac{L(M)}{M^2} f^B(q_0; T) \text{Im}\Pi_{EM}(M, q; \mu_B, T)$$

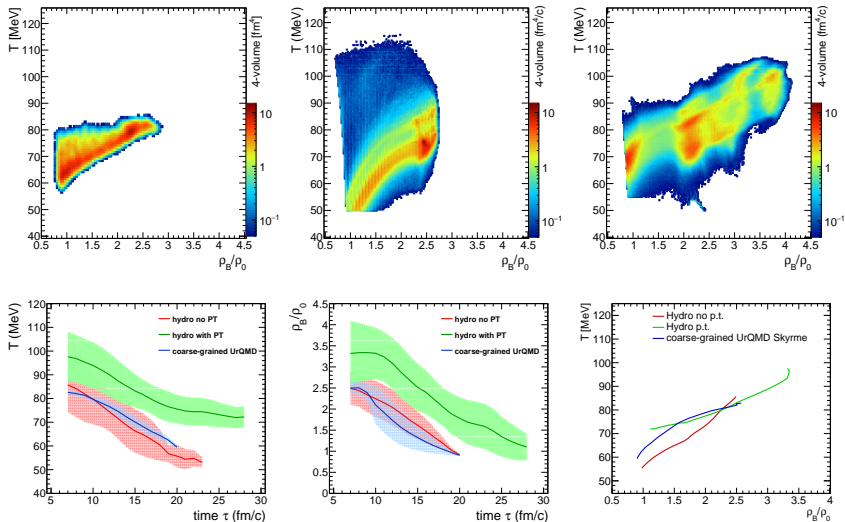
- M : virtual photon mass = dilepton invariant-mass, M_{ee} ($= \sqrt{q_0^2 - q^2}$)
- α_{EM} : electromagnetic coupling constant
- $\text{Im}\Pi_{EM}$: EM spectral-function of the QCD medium
- $f^B(q_0; T)$: thermal Bose distribution
- $L(M)$: lepton phase-space factor
- **High-Acceptance Di-Electron Spectrometer**
- SIS18 BES with Au+Au collisions at 1.23 AGeV measure M through M_{ee}
- Hadronic transport model, using UrQMD
- Hydrodynamic evolution, without first-order phase transition
- Hydrodynamic evolution, with first-order phase transition

Hydrodynamic simulations



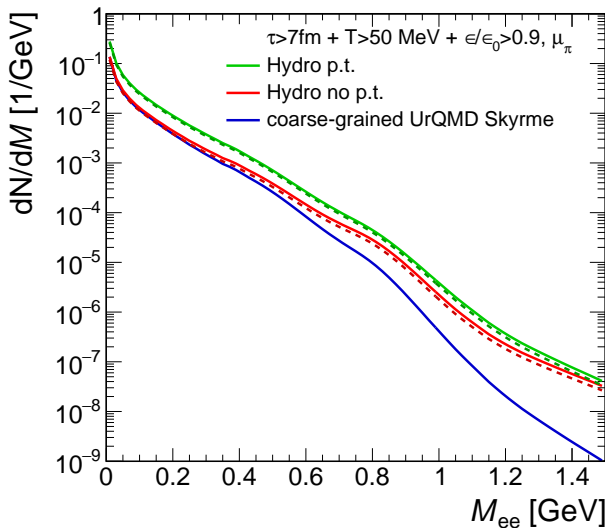
- Two different $Q\chi P$ EoS's as inputs
- Two hydro. sim.'s with three impact parameters – 2 fm, 4 fm & 7 fm
- T & ρ_B obtained as functions of x and $t \rightarrow \epsilon$ and M_{ee}

Bulk evolution



- Higher T & ρ_B and longer system-lifetimes in first-order scenario

Dilepton spectra



[F. Seck *et al*; Phys. Rev. C 106 (2022) 014904]

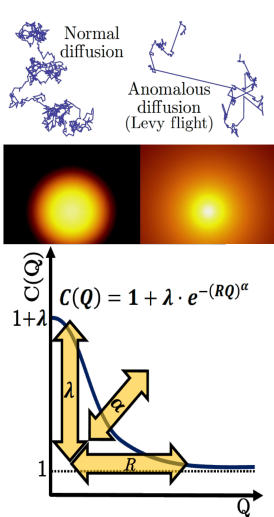
- FOPT doubles low-mass dilepton-yield; over that from crossover; due to prolonged system-lifetime caused by mixed-phase-formation

Outcomes & outlook

- Considerable influence of LG transition on cumulant values
- Acceptable nuclear-matter properties with stiff EoS (excluded volumes)
- Corroborated values for max. mass, canonical mass & canonical radius of NS's
- Double dilepton-yield with FOPT; w.r.t. crossover (prolonged system-lifetime)
- Hadronic re-scattering in dilute phases suppress effects of EoS-driven expansion
- Measurements of initial-state fluctuations, with effective theories
- Model simulations using non-zero net-strangeness and net-isospin
- Dynamic simulations of ultra-relativistic heavy-ion collisions, with UrQMD
- Application of obtained EoS to nuclear-astro. simulations of neutron stars.
- Exploration of similarities between NS-mergers and HIC's.
- Investigations into empirical observables as signatures of critical phenomena
- Probing of finite-size systems, with momentum cut-offs & Matsubara-sums
- Search for exotic phases of matter & hypernuclei, with effective theories
- Examination of magnetic field effects on phases of strongly-interacting matter

PART II: EXPERIMENT

Charged-kaon femtoscopy



- Mapping geometry of source \rightarrow momentum correlations of like-sign kaon-pairs:
 $C(q) = 1 + \tilde{D}(q)$; $\tilde{D}(q)$: FT of pair-source $D(r)$
- Usually assumed shape for $D(r)$ – Gaussian
- Generalization – Lévy distribution:
 $\mathcal{L}(r; \lambda, R) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-(RQ)^\alpha} e^{iQr} dQ$
- R : Lévy-scale, λ : correlation-strength, α : Lévy-exponent, Q : integration variable
- $\alpha = 2$: Gauss; $\alpha < 2$: power-law;
 $\alpha = 1$: Cauchy (or, exponential)
- Possible reasons for non-Gaussian sources:
 - Proximity to CEP: irrelevant at 200 GeV
 - Jet fragmentation: not possible in A+A
 - Anomalous diffusion: viable in A+A at 200 GeV

Hadronic re-scattering

- Evidence of non-Gaussian source-distribution for pions found in Au+Au collisions at PHENIX & STAR
- Extracted coordinate-space distributions show heavy tail
- Hydrodynamic calculations assume idealised freeze-out: sudden jump in mean-free-path from 0 to ∞
- More realistic scenario – hadronic re-scattering:
 - System cools & dilutes with expanding hadron-gas
 - Mean-free-path gradually diverges to ∞ , in finite time-interval
 - Re-scattering occurs in time-dependent mean-free-path-system
 - Anomalous diffusion – experimentally observed as power-law-shaped tails in coordinate-space distributions
 - In contrast to Gaussian, strongly-decaying tails for normal diffusion

- Momentum-space diffusion equation:

$$\frac{\partial W}{\partial t} = -K_n Q^2 W(Q, t)$$

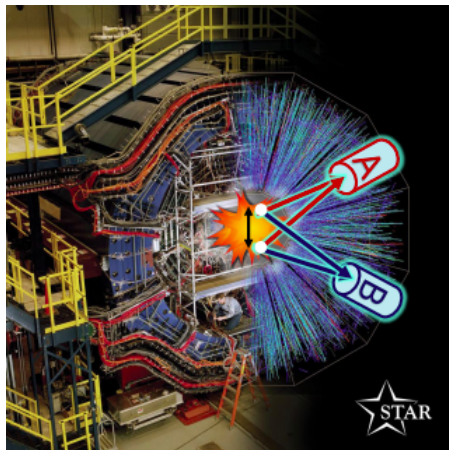
- K_n : normal diffusion constant
 - Q : momentum
 - t : time
 - $W(Q, t)$: momentum-space probability distribution
- Coordinate-space solution: $W(r, t) = \frac{1}{\sqrt{4\pi K_n t}} e^{-\frac{r^2}{4K_n t}} \rightarrow$ Gaussian

- Coordinate-space diffusion (generalised Fokker-Planck) equation:

[T. Csörgő, S. Hegyi, T. Novák & W. Zajc; AIP Conf. Proc. 828 (2006) 1, 525-532]

$$\frac{\partial W}{\partial t} + v \frac{\partial W}{\partial r} + \frac{F(r)}{m} \frac{\partial W}{\partial v} = \eta_{\alpha'} D_t^{1-\alpha'} L_{\text{FP}} W(r, v, t)$$

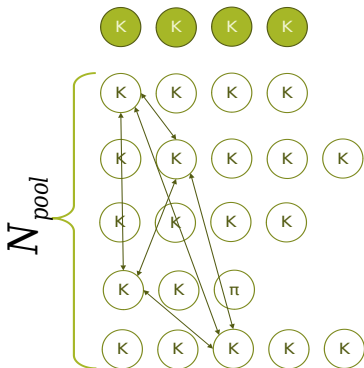
- Momentum-space solution: $W(Q, t) = e^{-tK^\alpha |Q|^\alpha}$
 - $W(Q, t)$: characteristic function (FT) of Lévy-stable source-distributions
 - α : Lévy-exponent
 - K : anomalous diffusion constant



- Solenoidal Tracker At RHIC
- Colliding ^{238}U , ^{197}Au , ^{63}Cu , ^{96}Zr , ^{96}Ru , ^{27}Al , ^3He , d & p
- Multiple centre-of-mass energies ($\sqrt{s_{\text{NN}}}$) for BES-I & BES-II
- Measurement: RHIC BES (2016) with Au+Au collisions at 200 GeV
- PID: dE/dx for K^+ , K^-

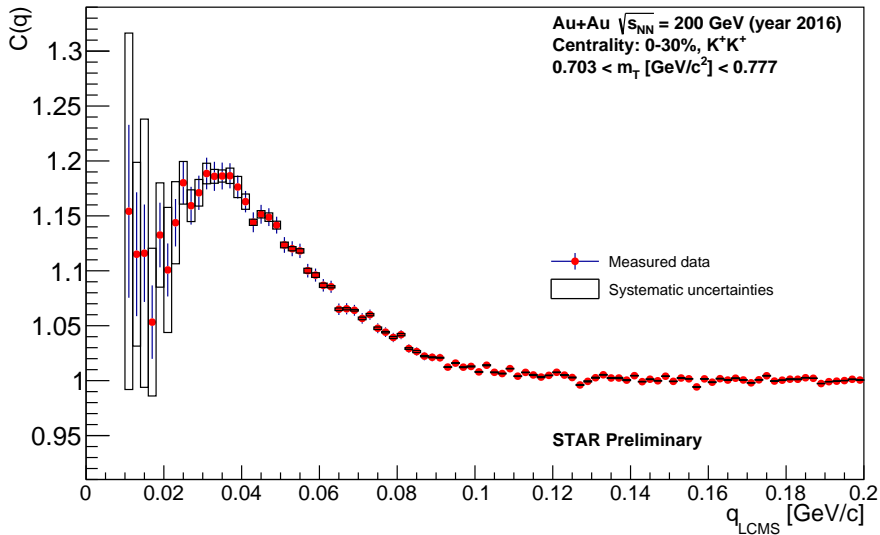
Event & track processing

- Event processing:
 - 3.06 billion events from 2016 RHIC beam-energy scan (BES) at 200 GeV in STAR's PicoDST file-storage
 - Trigger cuts (VPD, TPC, etc.) bring no. of events down to 2.59B
 - 0-30% centrality cut further reduces no. of events to 776 million
 - 52.8% of 776M events processed to get particle-tracks for analysis
- Track processing:
 - Tracks read in & cut (PID, N_{Hits} , etc.); $A(q)$ obtained
 - Pair cuts (FMH, SL & $\Delta z - \Delta u$) applied
 - Particles from current event stored in pool; events mixed
 - Over-weighting of events avoided \rightarrow only one particle selected from one event; $B(q)$ & $C(q)$ obtained
 - $C(q)$ fit with Coulomb-corrected Lévy-function
 - Fit parameters extracted & plotted with systematic uncertainties.



- Momentum (q) measured in Longitudinally Co-Moving System:
 $q_{\text{LCMS}} = |\vec{p}_1 - \vec{p}_2|_{\text{LCMS}}$
- Spherical symmetry in q_{LCMS} ideal for 1D analysis of 3D system
- $A(q)$ - kaon pairs from same event
- $B(q)$ - kaon pairs from mixed event
- Mixed event created by randomly selecting kaon-pairs from pool
- Correlation-function:
 $C(q) = A(q)/B(q)$
- 3 m_T bins used;
 $m_T = \sqrt{m^2 + (k_T/c)^2}$
- Lévy-type correlation function:
 $C(q) = 1 + \lambda \cdot e^{-(Rq)^\alpha}$

Correlation-function



- Correlation-function shows Bose-Einstein-peak & Coulomb-hole

Lévy-function & Coulomb-correction

- Bowler-Sinyukov formula with Coulomb-repulsion:

[Y. Sinyukov et al; Phys. Lett. B 432 (1998) 248-257]

$$C(q) = \left[1 - \lambda + \lambda \cdot K(q) \cdot \left(1 + e^{-(Rq)^\alpha} \right) \right] \cdot N \cdot (1 + \varepsilon q) ,$$

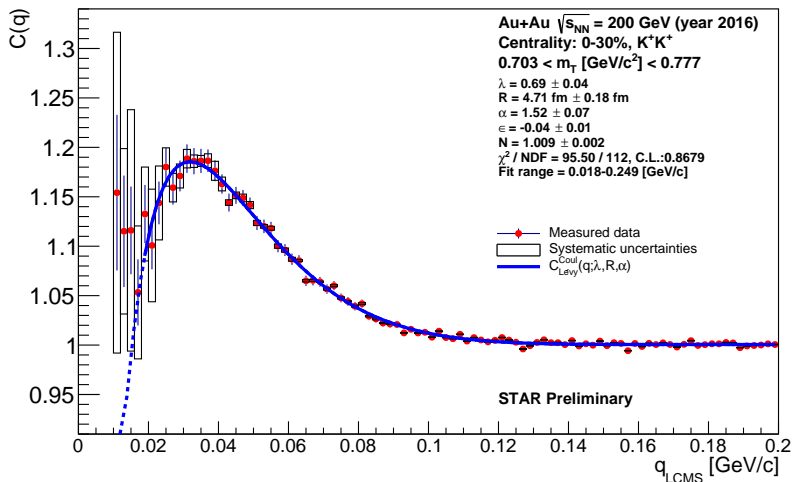
- $N \cdot (1 + \varepsilon q)$: assumed linear background
- Coulomb-correction:

[M. Csanád, S. Lökös & M. Nagy; Phys. Part. Nucl. 51 (2020) 3, 238-242]

$$K(q; \alpha, R) = \frac{\int D(r) |\psi^{\text{Coul}}(r)|^2 dr}{\int D(r) |\psi^0(r)|^2 dr} ,$$

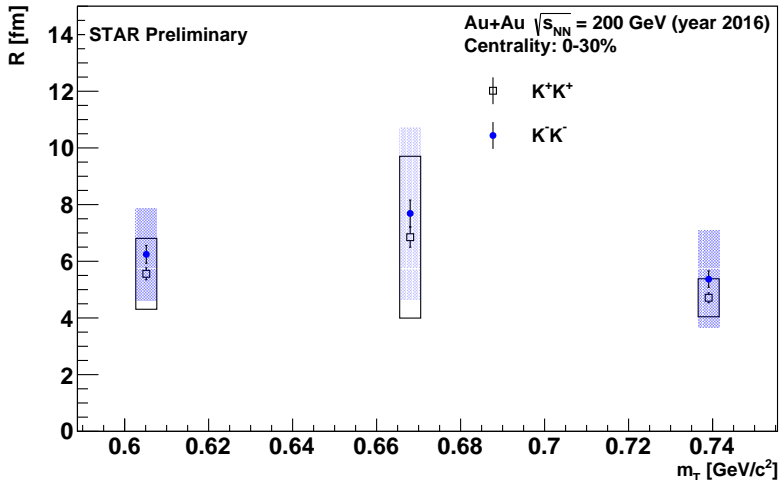
- $D(r)$: spatial pair-distribution
- ψ^0 : 2-particle plane-wave
- ψ^{Coul} : Coulomb-wave
- $K(q; \alpha, R)$ modified for kaons & calculated numerically

Sample fit



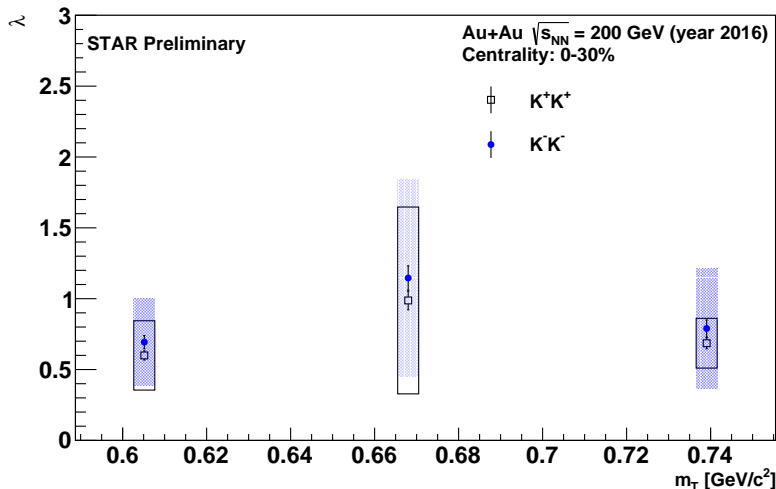
- Measured $C(q)$ agrees quantitatively with best fit over entire q -range
- $N \approx 1$ & $\epsilon \approx 0$ from fitting – linear contribution negligible

Lévy-scale R



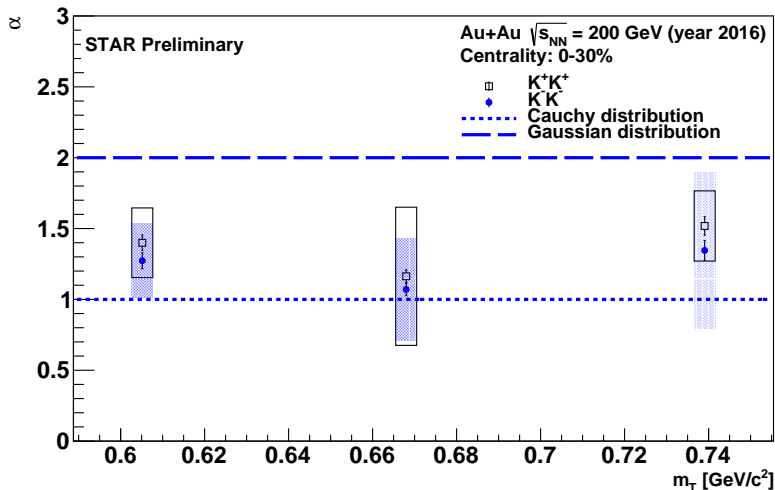
- Kaon-homogeneity length: very weak dependence on m_T ; large uncertainties
- Possible slight decrease; not contradicting hydro.-predictions
- Similar to PHENIX pion data: $R_\pi(m_T=0.6-0.7 \text{ GeV}/c^2) \approx 5-7 \text{ fm}$
[PHENIX Collaboration; Phys. Rev. C 97 (2018) 6, 064911]

Correlation-strength λ



- Intercept of correlation-function – Core-Halo model: $\lambda = N_C / (N_C + N_H)$
[T. Csörgő, B. Lorstad & J. Zimányi; Z. Phys. C 71 (1996) 491-497]
- Close to unity; in line with expected, small fraction of decay-kaons

Lévy-exponent α



- May describe extent of anomalous diffusion
- $\alpha \approx 1.0 - 1.5$ for kaons, similar to PHENIX pion results: $\alpha_\pi \approx 1.2$
[PHENIX Collaboration; Phys. Rev. C 97 (2018) 6, 064911]
- Suggests non-Gaussian source-shape for charged kaons, similar to pions

Outcomes & outlook

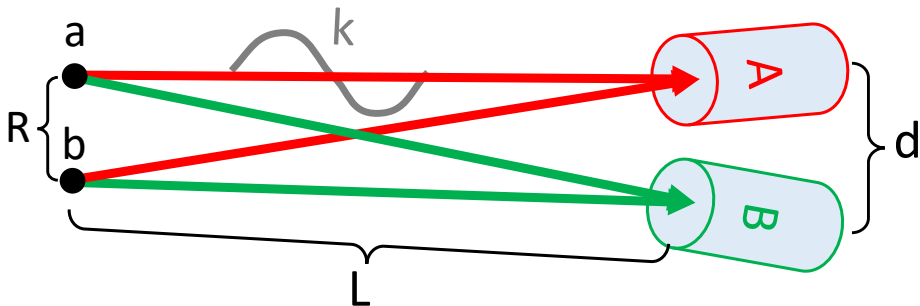
- Preliminary analysis suggests a non-Gaussian source-shape for pairs of charged kaons produced in heavy-ion collisions
- Lévy-stability-exponent α is comparable to that of pions
- Anomalous diffusion is not the sole reason for the heavy tails; since $\alpha_K \approx \alpha_\pi$, instead of $\alpha_K \not\approx \alpha_\pi$
- Full systematic-uncertainty-analysis is required to achieve definitive conclusions about the geometry of the source
- Similar measurements at lower centre-of-mass energies would be interesting as probes for the critical end-point

Thank you for your attention!

অন্তরে অতৃপ্তি র'বে সাঙ্গ করি' মনে হবে
শেষ হয়ে হইল না শেষ।

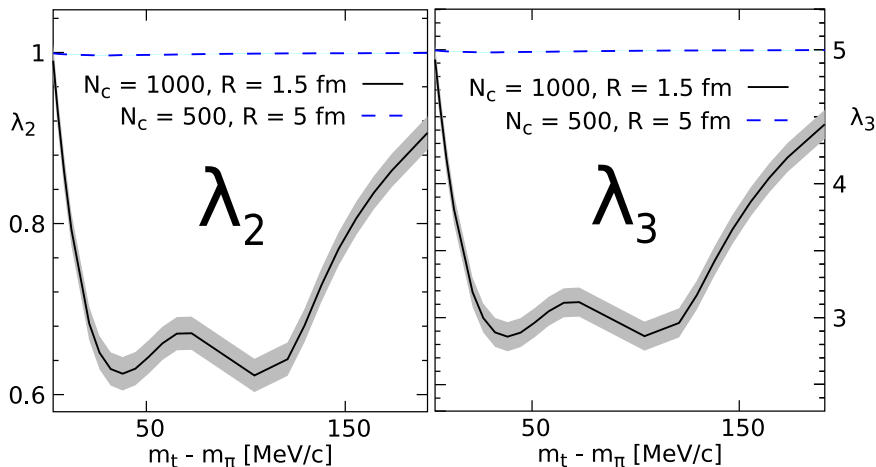
APPENDIX

Random phases



- a and b as sources, A and B as detectors
- R and d as distance between the sources and detectors, respectively
- k as the phase difference and L as the path length
- Two- and three-particle correlation-strengths, with random-phase:
 - $\lambda_2 = C_2(0) - 1 = e^{-2\sigma_\phi^2}$
 - $\lambda_3 = C_3(0) - 1 = 3e^{-2\sigma_\phi^2} + 2e^{-3\sigma_\phi^2}$

Correlation-strengths



[M. Csanad, A. Jakovac, S. Lokos, A. Mukherjee & S. K. Tripathy; Gribov-90 Memorial Volume 7 (2021) 261–273]

- Low- m_t decrease of $\lambda_{2,3}$
- Magnitude strongly depends on charge density